

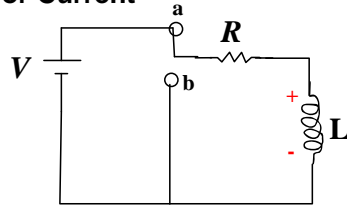
PHYS102  
DC-Circuits  
with  
Inductors

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April 13, 2007

LR Circuit - Equation for Current . . . . .	2
Inductive Time Constant . . . . .	3
Current Behavior - Graphing . . . . .	4
Graphing - EMF of Inductor . . . . .	5
Disconnecting the Battery . . . . .	6
Current Decay through LR Circuits . . . . .	7
LR - RC Circuit Comparison . . . . .	8
<b>Energy</b>	<b>9</b>
Energy Stored in Inductors . . . . .	9
Current and Energy in Inductors . . . . .	10
Magnetic Energy Density . . . . .	11
<b>LC - Oscillator</b>	<b>12</b>
Energy Oscillations . . . . .	12
LC - Circuit . . . . .	13
Charge Oscillations . . . . .	14

### LR Circuit - Equation for Current



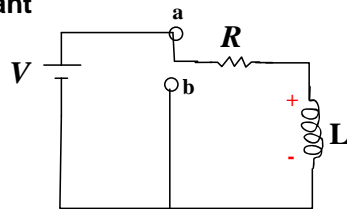
$$\Rightarrow \int \frac{dI}{(I - V/R)} = - \int \frac{R dt}{L}$$

$$I = \frac{V}{R} (1 - e^{-Rt/L})$$

PHYS102 NOTE:  $L/R$  has units of seconds.  $L/R$  has units of seconds.

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### Inductive Time Constant



- Continuation from last slide.

$$I = \frac{V}{R} (1 - e^{-Rt/L})$$

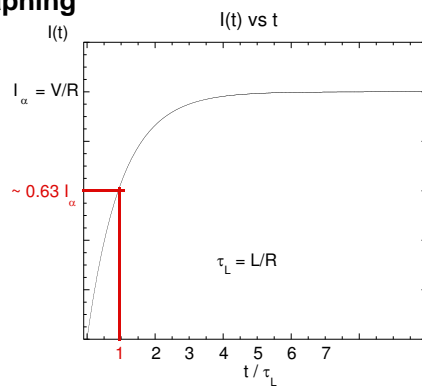
$$I = \frac{V}{R} (1 - e^{-t/\tau_L}) \quad (\text{where } \tau_L \equiv L/R)$$

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$\tau_L$  is called the “inductive” time constant for the circuit. Circuits with Inductors – slide 3

## 0.1 Graphing Current and EMF

### Current Behavior - Graphing



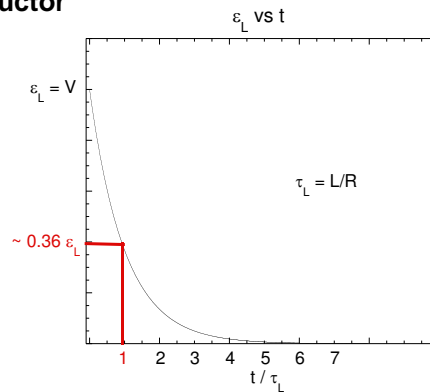
$$I(t) = \frac{V}{R} \left( 1 - e^{-Rt/L} \right)$$

- The current through the inductor builds up over time (just like we stated conceptually).

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- What happens to the EMF in the inductor?

### Graphing - EMF of Inductor



$$I(t) = \frac{V}{R} \left( 1 - e^{-Rt/L} \right)$$

$$\varepsilon_L = -L \frac{dI}{dt} = V e^{-Rt/L}$$

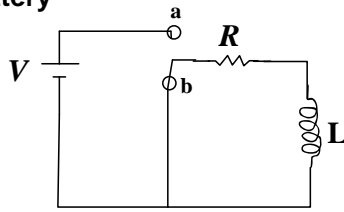
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- The EMF in the inductor approaches zero as the current in the circuit reaches equilibrium (i.e., current does not fluctuate).

## 0.2 Removing the Battery

### Disconnecting the Battery



- After a very long time, the switch is thrown into position (b).
  - The battery is disconnected from the rest of the circuit.
- Writing Kirchhoff's Rules for this loop:

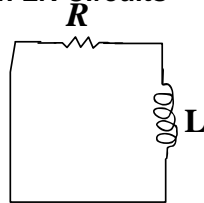
$$-I R - |\varepsilon_L| = 0.$$

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$$\Rightarrow L \frac{dI}{dt} = -I R$$

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### Current Decay through LR Circuits



- Continuation from the last slide.

$$L \frac{dI}{dt} = -I R$$

$$\Rightarrow \frac{dI}{I} = -\frac{R}{L} dt$$

$$\Rightarrow I(t) = I_0 e^{-t/\tau_L}$$

PHYS102 presents the current through the inductor right before the switch was thrown into position (b).

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### LR - RC Circuit Comparison

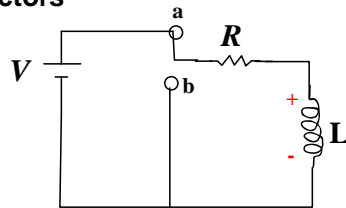
- The equations used to describe the current through the inductor (RL circuit) and the charge on a capacitor (RC circuit) look identical in form.
- We calculated the energy stored in a capacitor by considering the buildup of charge.
  - What is the energy of an inductor with current flowing through it?
- Start with the circuit first discussed in this lecture!

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Circuits with Inductors – slide 8

## 0.3 Power

## Energy Stored in Inductors



- From Kirchhoff's Rules:

$$\Rightarrow V = IR + L \frac{dI}{dt}$$

- Multiply both sides of the equation by  $I$ .

$$\Rightarrow VI = I^2 R + LI \frac{dI}{dt}$$

- The  $VI$  term is the power delivered by the battery. The  $I^2 R$  term is the power dissipated by the resistor.

This means the  $LI \frac{dI}{dt}$  term is the power stored in the inductor.

Circuits with Inductors – slide 9

## 0.4 Power in Inductors

## Current and Energy in Inductors

- Power is the rate of change of Energy per time.

$$P_L = LI \frac{dI}{dt} = \frac{dU}{dt}$$

$$\Rightarrow dU_L = LI dI$$

$$\Rightarrow U_L = \int LI dI = \frac{1}{2} LI^2$$

- The energy stored in the inductor varies as  $I^2$  (much like  $Q^2$  for a capacitor).
- Is there a general expression that does not depend on the self-inductance? In other words, is there a relationship between the amount of energy and the magnetic field?

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## 0.5 Energy and B-fields

### Magnetic Energy Density

- Consider a solenoid with radius  $a$ , length  $l$ , and  $n$  number of turns per unit length.

$$B = \mu_0 n I$$

$$L = n^2 \mu_0 l (\pi a^2) = n^2 \mu_0 (Al)$$

$$\Rightarrow U_B = \frac{1}{2} L I^2 = \frac{1}{2} n^2 \mu_0 (Al) \left( \frac{B}{\mu_0 n} \right)^2$$

$$\Rightarrow U_B = \frac{1}{2 \mu_0} B^2 (Al)$$

- $(Al)$  is the volume where the magnetic field exists. Think energy per volume.

$$u_B = U_B/V = \frac{B^2}{2 \mu_0}$$

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## LC - Oscillator

slide 12

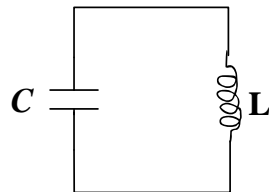
### Energy Oscillations

- Since capacitors store electrical energy and inductors store magnetic energy, we could place a fully charged capacitor in series with an inductor.
- The electrical energy should be transferred to magnetic energy, and then the magnetic energy should get transferred back into electrical energy. This cycle should repeat itself. Let's prove it.
- Since energy stored in a capacitor is proportional to  $Q^2$ , it suffices to prove that the charge on the capacitor "oscillate".

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Circuits with Inductors – slide 12

### LC - Circuit



- Start with a *fully charged capacitor* and place it in series with an inductor as shown above.
- Write down the total energy of the system at some time  $t$  after the capacitor is connected to the inductor.

$$U_T = \frac{1}{2} \frac{Q^2}{C} + \frac{1}{2} L I^2$$

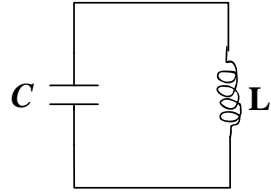
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$$\frac{dU_T}{dt} = 0 = \frac{Q}{C} \frac{dQ}{dt} + L I \frac{dI}{dt}$$

$$0 = \frac{Q}{C} + L \frac{d^2 Q}{dt^2}$$

## Charge Oscillations



$$\frac{d^2 Q}{dt^2} = -Q/LC = -\omega^2 Q$$

- This equation defines simple harmonic motion with an angular frequency  $\omega = \frac{1}{\sqrt{LC}}$ . The charge on the capacitor,  $Q$ , is undergoing simple harmonic motion.

► **PHYS 102** mechanics, we know a solution for  $Q(t)$ :

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$$Q(t) = Q_0 \sin(\omega t + \phi)$$